Dynamics of strong and radiative decays of D_s-mesons in the hadrogenesis conjecture

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Abstract The positive parity scalar $D_{s0}^*(2317)$ and axial-vector $D_{s1}^*(2460)$ charmed strange mesons are generated by coupled-channel dynamics through the s-wave scattering of Goldstone bosons off the pseudoscalar and vector $D(D_s)$-meson ground states. The specific masses of these states are obtained as a consequence of the attraction arising from the Weinberg-Tomozawa interaction in the chiral Lagrangian. Chiral corrections to order Q_χ^2 are calculated and found to be small. The $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons decay either strongly into the isospin-violating $\pi^0 D_s$ and $\pi^0 D_s^*$ channels or electromagnetically. We show that the π^0-η and $(K^0 D^+ - K^+ D^0)$ mixings act constructively to generate strong widths of the order of 140 keV and emphasize the sensitivity of this value to the KD component of the states. The one-loop contribution to the radiative decay amplitudes of scalar and axial-vector states is calculated using the electromagnetic Lagrangian to chiral order Q_χ^2. We show the importance of taking into account processes involving light vector mesons explicitly in the dynamics of electromagnetic decays. The radiative width are sensitive to both ηD_s and KD components, hence providing information complementary to the strong widths on the positive parity D_s-meson structure.

Key words charmed mesons, $D_{s0}^*(2317)$, $D_{s1}^*(2460)$

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1 Introduction

The positive parity scalar $D_{s0}^*(2317)$ and axial-vector $D_{s1}^*(2460)$ charmed strange mesons have peculiar properties. Their masses are significantly lower than expected in quark models [1]. Their strong decays into the $\pi^0 D_s$ and $\pi^0 D_s^*$ channels violate isospin and their radiative decays appear quite selective. We recall the corresponding data.

The present upper limit for the $D_{s0}^*(2317)$ total width is $\Gamma < 3.8$ MeV [2]. Its radiative width should be extremely small according to the constraint [3]

$$\frac{\Gamma\left[D_{s0}^*(2317) \to D_s^*(2112)\gamma\right]}{\Gamma\left[D_{s0}^*(2317) \to D_s(1968)\pi^0\right]} < 0.059. \qquad (1)$$

The $D_{s1}^*(2460)$-meson's total width is less than 3.5 MeV [2]. Constraints on its radiative decays are as follows [3],

$$\frac{\Gamma\left[D_{s1}^*(2460) \to D_s(1968)\gamma\right]}{\Gamma\left[D_{s1}^*(2460) \to D_s^*(2112)\pi^0\right]} = 0.31 \pm 0.06, \qquad (2)$$

$$\frac{\Gamma\left[D_{s1}^*(2460) \to D_s^*(2112)\gamma\right]}{\Gamma\left[D_{s1}^*(2460) \to D_s^*(2112)\pi^0\right]} < 0.16, \qquad (3)$$

$$\frac{\Gamma\left[D_{s1}^*(2460) \to D_{s0}^*(2317)\gamma\right]}{\Gamma\left[D_{s1}^*(2460) \to D_s^*(2112)\pi^0\right]} < 0.22. \qquad (4)$$

The purpose of the work [4] summarized in this paper is to provide a consistent explanation of these properties of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons in the hadrogenesis conjecture [5, 6] extending earlier investigations [7–9] restricted to the study of the masses of these states. An important finding of our analysis is that masses, strong and radiative decays provide different and complementary constraints on the dynamics underlying the generation of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons and hence should be described together in the same theoretical framework.

We recall briefly in Section 2 the main assumptions underlying the hadrogenesis conjecture. We outline in Sections 3 how the strong and electromagnetic decays of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons

are calculated and mention the dynamical origin of the features summarized above. We refer to Ref. [4] for a complete presentation of the formalism and of the numerical results.

2   Hadrogenesis

The hadrogenesis conjecture [5, 6] applied to heavy-light meson states [8, 9] has been successful in generating hadron resonances by coupled-channel dynamics with $J^\pi = 0^+, 1^+$ quantum numbers. For D_s-mesons, they are produced by the s-wave scattering of Goldstone bosons off D-meson ground state triplets. The D-meson ground state triplet is composed either of the pseudoscalars $\{D^0(1864), -D^+(1869), D_s^+(1968)\}$ to build the 0^+ state or of the vectors $\{D_\mu^{*0}(1864), -D_\mu^{*+}(1869), D_{s\mu}^{*+}(1968)\}$ to generate the 1^+ state. For the $D_{s0}^*(2317)^+$ meson, the calculation involves the ηD_s^+, $K^0 D^+$ and $K^+ D^0$ channels coupled further to the $\pi^0 D_s^+$ channel through isospin-mixing parameters. Analogously we consider for the $D_{s1}^*(2460)^+$ meson the ηD_s^{*+}, $K^0 D^{*+}$, $K^+ D^{*0}$ and $\pi^0 D_s^{*+}$ channels. Both the $D_{s0}^*(2317)^+$ and $D_{s1}^*(2460)^+$ mesons are $SU(3)$ anti-triplet bound states.

The interaction between the Goldstone bosons and the scalar D-meson triplet is governed by the leading order chiral Lagrangian density [8, 9]

$$\mathcal{L} = \frac{1}{4}\,\mathrm{tr}\,(\partial_\mu \Phi)\,(\partial^\mu \Phi) - \frac{1}{4}\,\mathrm{tr}\,\chi_0\,\Phi^2 + (\partial_\mu D)\,(\partial^\mu \bar{D}) -$$
$$D\,M_{0^-}^2\,\bar{D} + \frac{1}{8\,f^2}\left\{(\partial^\mu D)\,[\Phi,(\partial_\mu \Phi)]_-\,\bar{D} - D\,[\Phi,(\partial_\mu \Phi)]_-\,(\partial^\mu \bar{D})\right\}, \tag{5}$$

where Φ is the pseudoscalar octet

$$\Phi = \begin{pmatrix} \pi^0 + \dfrac{1}{\sqrt{3}}\eta & \sqrt{2}\pi^+ & \sqrt{2}K^+ \\ \sqrt{2}\pi^- & -\pi^0 + \dfrac{1}{\sqrt{3}}\eta & \sqrt{2}K^0 \\ \sqrt{2}K^- & \sqrt{2}\bar{K}^0 & -\dfrac{2}{\sqrt{3}}\eta \end{pmatrix}. \tag{6}$$

and D the triplet field. We use the notation $\bar{D} = D^\dagger$. The parameter f in (5) is the octet meson decay constant. It defines the scale of chiral symmetry breaking and is approximatively known from the weak decay of the charged pions. A precise determination of f requires a chiral $SU(3)$ extrapolation of some data set. We will use the value $f = 90$ MeV obtained from a detailed study of pion- and kaon-nucleon scattering data in Ref. [6]. This parameter determines the leading s-wave interaction of the Goldstone bosons with the open-charm meson fields. The ground-state

scalar D-meson mass matrix is denoted by M_{0^-}. The mass term of the Goldstone bosons is proportional to the quark-mass matrix χ_0.

An expression similar to (5) is derived for the interaction between the Goldstone bosons and the vector D-meson triplet,

$$\mathcal{L} = -(\partial_\mu D^{\mu\alpha})\,(\partial^\nu \bar{D}_{\nu\alpha}) + \frac{1}{2}\,D^{\mu\alpha}\,M_{1^-}^2\,\bar{D}_{\mu\alpha} -$$
$$\frac{1}{8\,f^2}\left\{(\partial^\nu D_{\nu\alpha})\,[\Phi,(\partial_\mu \Phi)]_-\,\bar{D}^{\mu\alpha} - D_{\nu\alpha}\,[\Phi,(\partial_\nu \Phi)]_-\,(\partial_\mu \bar{D}^{\mu\alpha})\right\}, \tag{7}$$

where we represent the 1^- D-mesons in terms of antisymmetric tensor fields. This particular representation has the advantage of leading to gauge-invariant expressions for the radiative processes considered in this work. The results obtained in Refs. [5] and [6] with the vector representation were reformulated in the tensor representation [4]. Using these chiral Lagrangians at leading order, strong attraction was found in the chiral $SU(3)$ antitriplet states to which the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ belong. The key parameter underlying this result is the chiral symmetry breaking scale f. To improve on the leading order calculation, chiral correction terms were included in a systematic way to order Q_χ^2 [9]. They take into account the s- and u-channel exchanges of D-mesons and local counter terms. The corresponding vertices introduce additional coupling constants constrained by data when available, heavy quark symmetry and the large N_c limit of QCD [4].

We note in addition that a striking prediction of the scenario outlined above is a clear measurable signal of an exotic axial state in the ηD^* invariant mass distribution with a mass of 2568 MeV [4]. It lies above the ηD^* threshold and has a width of about 18 MeV. It could be discovered by ongoing experiments once the invariant ηD^* mass distribution is analyzed.

3   Strong and radiative decays of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$

The strong decays of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ into the $\pi^0 D_s$ and $\pi^0 D_s^*$ respectively violate the isospin symmetry. The origin of these processes is the difference between the up- and down-quark masses which induces two kinds of isospin mixings: the π^0-η mixing characterized by an angle ϵ (determined to be 0.01 [10]) and the $(K^0 D^+$-$K^+ D^0)$ mixing through $I{=}1$ transitions. These effects act constructively [4]. The isospin-violating strong widths of the $D_{s0}^*(2317)$ and

$D_{s1}^*(2460)$ mesons are sensitive to both the coupled-channel dynamics and the angle ϵ. We find that the strong widths of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ are comparable (as suggested by the heavy-quark symmetry of QCD) and of the order of 140 keV [4]. We emphasize that these widths are largely dominated by the terms arising from the $(K^0 D^+ - K^+ D^0)$ and $(K^0 D^{*+} - K^+ D^{*0})$ mixings respectively. The strong widths of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons are therefore good measures of their KD and KD* structures.

Our calculated value of 140 keV is smaller than the experimental upper limits by a factor of the order of 20. More accurate determinations of these widths would clearly be very useful in unraveling the dynamics underlying these decays.

To calculate the radiative decays of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons we define the Lagrangian for electromagnetic interactions to chiral order Q_χ^2. It involves four contributions. First, we gauge the Lagrangian describing strong interactions. Only the 3-point vertices involved in the chiral correction terms

matter. Secondly, we consider terms of chiral order Q_χ^2 proportional to the electromagnetic tensor $F_{\mu\nu}$ and corresponding to 4-point vertices ($D^*D\pi\gamma$ for example). A third kind of terms of chiral order Q_χ^2 and proportional to $F_{\mu\nu}$ describes the anomalous processes induced by 3-point vertices ($D^*D\gamma$) and the magnetic moments of the charmed vector mesons. A fourth set of terms is introduced to investigate the role of light vector mesons in the radiative transitions.

We display in Fig. 1 the diagrams contributing to the electromagnetic decay of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons in the hadrogenesis picture. The evaluation of these graphs is described in Ref. [4]. Our numerical results are linked to specific choices of parameters. The explicit introduction of light vector mesons in the radiative decay dynamics in needed to describe the available data summarized in Eqs. (1)–(4) with natural size parameters. The light vector mesons play a role through gauge-invariant 3- and 4-point vertices proportional to the electromagnetic field strength tensor.
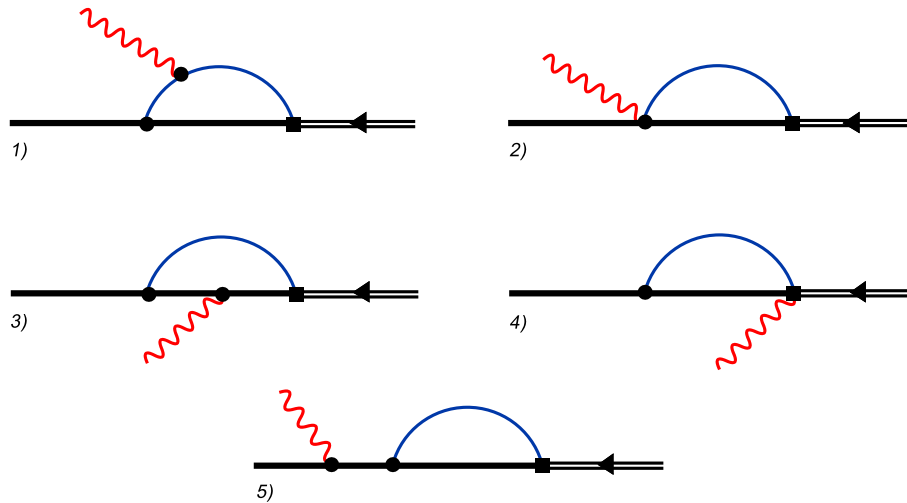

Fig. 1.   (color online).   Diagrams contributing to the decay amplitude of a scalar or axial vector molecule. Solid lines represent the propagation of the pseudo-scalar or vector mesons, where the thin lines stand for the light mesons and the thick lines for the heavy mesons. The wavy line is the photon.

Specific dynamical effects emerge from our work. The full relativistic loop computations induce significant deviations from heavy-quark symmetry as expected from the semi-heavy character of the charm-quark mass. The ηD_s and ηD_s^* channels play an important role in radiative decays. These decays are therefore important observables to understand the coupled-channel structure of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons and should be discussed together with the strong decays which are mainly sensitive to

the KD and KD* components of the states.

Because of large cancellations between graphs involving ηD_s and KD channels, we are not able to make definite predictions at this point. Taking natural sets of parameters chosen to reproduce the ratio (2), we obtain for the ratios (1), (3) and (4) respectively the ranges of values 0.014–0.046, 0.08–0.15 and 0.001–0.004. These numbers should be taken as mere indications and incentives to measure more precisely these experimental ratios. The investigation of the

role of vector mesons indicates that the K^*D^* and ϕD_s^* channels could contribute significantly to the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ radiative decays.

More accurate data, further coupled-channel studies and an improved knowledge of the couplings involved in the calculation (through measurements or lattice studies) are needed to make progress in the dynamical understanding of the structure of the $D_{s0}^*(2317)$ and $D_{s1}^*(2460)$ mesons.

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